Nanodiamond-Quantum Sensors Reveal Temperature Variation Associated to Hippocampal Neurons Firing

Original
Nanodiamond-Quantum Sensors Reveal Temperature Variation Associated to Hippocampal Neurons Firing / Petrini, Giulia; Tomagra, Giulia; Bernardi, Ettore; Moreva, Ekaterina; Traina, Paolo; Marcantoni, Andrea; Picollo, Federico; Kvaková, Klaudia; Cígler, Petr; Degiovanni, Ivo Pietro; Carabelli, Valentina; Genovese, Marco. - In: ADVANCED SCIENCE. - ISSN 2198-3844. - 9:28(2022), pp. 1-11. [10.1002/advs.202202014]

Availability:
This version is available at: 11696/76179 since: 2023-02-28T14:23:46Z

Publisher:
WILEY

Published
DOI:10.1002/advs.202202014

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Temperature is one of the most relevant parameters for the regulation of intracellular processes. Measuring localized subcellular temperature gradients is fundamental for a deeper understanding of cell function, such as the genesis of action potentials, and cell metabolism. Notwithstanding several proposed techniques, at the moment detection of temperature fluctuations at the subcellular level still represents an ongoing challenge. Here, for the first time, temperature variations (1 °C) associated with potentiation and inhibition of neuronal firing is detected, by exploiting a nanoscale thermometer based on optically detected magnetic resonance in nanodiamonds. The results demonstrate that nitrogen-vacancy centers in nanodiamonds provide a tool for assessing various levels of neuronal spiking activity, since they are suitable for monitoring different temperature variations, respectively, associated with the spontaneous firing of hippocampal neurons, the disinhibition of GABAergic transmission and the silencing of the network. Conjugated with the high sensitivity of this technique (in perspective sensitive to < 0.1 °C variations), nanodiamonds pave the way to a systematic study of the generation of localized temperature gradients under physiological and pathological conditions. Furthermore, they prompt further studies explaining in detail the physiological mechanism originating this effect.

1. Introduction

On the one side, temperature regulates the speed of ion channel opening,[1] the pattern activity of a firing neuron,[2] and the vesicular dynamics at the presynaptic terminal,[3] on the other side, intracellular temperature is affected by a variety of biochemical reactions occurring during cell activity. Pioneering findings dating back to the late 70s,[4,5] associated temperature increases to the impulse propagation in nonmyelinated fibers of the olfactory nerve and, more recently, a theoretical explanation about the heat production and absorption by neurons during nervous conduction has been formulated.[6] Intracellular temperature variations have been probed to detect the phases of cell-cycle division[7] and the mitochondrial activity.[8] In the brain, temperature fluctuations are likely associated with changes in neuronal functions such as the genesis of action potentials, secretion at the level of synaptic terminals, or transmembrane ion...
transports\textsuperscript{[9–11]} Variation of intracellular temperature is also related to changes in cell metabolism, as observed through positron emission tomography and functional magnetic resonance imaging (fMRI).\textsuperscript{[12]}

Besides its fundamental role in cell physiology, the temperature may be altered under pathological conditions, such as in cancerous cells\textsuperscript{[13]} which display higher metabolic activity, in neurodegeneration, such as Parkinson’s or Alzheimer’s diseases, where the process of abnormal protein aggregation is temperature-dependent,\textsuperscript{[14]} or during the onset of malignant hyperthermia, a pathology that causes rapid temperature increases following excessive muscle contraction.\textsuperscript{[15]}

Furthermore, an increase in local temperature, up to a few Celsius degrees, has been detected as triggered by laser heating,\textsuperscript{[16]} calcium stress,\textsuperscript{[17]} direct electric stimulation,\textsuperscript{[18]} or using drugs that increase the heat produced during cellular respiration.\textsuperscript{[16,18,19]} These results, in particular the ones related to the highest temperature variations, stimulated a debate since the power needed to justify this temperature growth has been calculated to exceed a few orders of magnitude what is expected by thermodynamic considerations in a model where thermal diffusion is characterized by conductive regime.\textsuperscript{[20]} Nonetheless, other authors\textsuperscript{[21,22]} suggested that, by taking into account the inhomogeneity of the cells, this gap can be strongly reduced.

However, detection of temperature fluctuations at the subcellular level still represents an ongoing challenge and several intracellular thermometry techniques are reported in the literature, ranging from fluorescent molecular thermometers,\textsuperscript{[23]} quantum dots,\textsuperscript{[24]} or rare-earth nanoparticles.\textsuperscript{[25]} Compared to the above-mentioned methods, nanodiamonds (ND) show a better biocompatibility,\textsuperscript{[26–27]} insensitivity to biological environment,\textsuperscript{[28]} more stable photoluminescence (PL) and a lower noise floor.\textsuperscript{[29]}

All-optical temperature measurement methods, based on the temperature-dependent shift of the luminescence spectra of color centers in diamond, such as silicon-vacancy (SiV), germanium vacancy (GeV), and tin vacancy (SnV) with strong zero-phonon line (ZPL), can be used.\textsuperscript{[30]} These methods are promising, but for the moment, apart from one very invasive experiment,\textsuperscript{[31]} there are few reported biosensing applications. For instance, in a notable experiment,\textsuperscript{[32]} gold nanorod–fluorescent ND hybrids were used as nanoheater/nanothermometer for studying living human embryonic kidney cells. In this case the NDs registered the temperature variation associated with the heating obtained by local laser irradiation of the nanorods. In another recent experiment,\textsuperscript{[33]} SiV in NDs were used to for imaging and temperature measurement in HeLa cells, even if only the average temperature of the bath (by controlling the incubator temperature) was measured, without single-cell resolution. It must be stressed anyway that, in both the above mentioned works, the temperature variation was induced by the experimenters (i.e., via laser irradiation or by changing the incubator temperature) and no authentic biological phenomenon was measured. Notwithstanding the undoubted promising perspective of all-optical methods, currently the most promising approach exploits nitrogen-vacancy (NV) centers.\textsuperscript{[29]}

For these reasons, NV color centers in ND combined with optically detected magnetic resonance (ODMR) technique can assume a dominating position in thermometry\textsuperscript{[34]} for biological application.\textsuperscript{[35,36]} Nevertheless, ODMR measurements protocols are based on microwave irradiation of the samples, which makes this technique more elaborated and requires awareness on the exploitable optical and microwave power providing the required sensitivity without damaging the cells.

ND was used for the first time in in vitro\textsuperscript{[37]} measurements of (of a heated gold nanoparticle inside a cell). Subsequently, an in vivo experiment was performed inside Caenorhabditis elegans worms, reaching a sensitivity as low as 1.4 °C Hz\textsuperscript{−1/2}. In the same experiment, thermogenic responses have been monitored during the chemical stimuli of mitochondrial uncouplers.\textsuperscript{[16]} Temperature gradients have been mapped at the subcellular level into a single human embryonic fibroblast\textsuperscript{[7]}; intracellular temperature mapping has been performed also in cultured primary cortical neurons, employing micro electrode arrays (MEAs) to demonstrate that the presence of NDs in primary cortical neurons does not elicit a neurotoxic response,\textsuperscript{[38]} in good agreement with our previous findings.\textsuperscript{[19]} Finally, heterogeneous temperature variations have been coupled with Ca\textsuperscript{2+} increases in HeLa cells.\textsuperscript{[19]}

Here, for the first time, we quantify the direct correlation of intracellular temperature variations with the modulation of neuronal activity. We demonstrate that sensors based on NV centers in NDs, interrogated via ODMR techniques, consisting of simultaneously applying a variable frequency microwave field and nonresonant laser radiation to the sample while acquiring the PL intensity (photon counts) emitted by a single ND.

The ODMR protocol used is based on the possibility to optically detect the negatively charged NV− center’s spin state and optically pump it into the |m_s = 0> sublevel. The mechanism is described as follows.

2. Results

In the presented experiments, temperature variations, associated to different patterns of neuronal activity, were detected using an ODMR measurement protocol, consisting of simultaneously applying a variable frequency microwave field and nonresonant laser radiation to the sample while acquiring the PL intensity (photon counts) emitted by a single ND.

The ODMR protocol used is based on the possibility to optically detect the negatively charged NV− center’s spin state and optically pump it into the |m_s = 0> sublevel. The mechanism is described as follows.

2.1. Optically Detected Magnetic Resonance (ODMR)

Depending on the NV electron spin state, two different decaying paths are possible. If the system is in the state |m_s = 0>, then a spin preserving transition occurs with single-photon emission at wavelength $\lambda = 637$ nm (ZPL). This transition is cyclic and one detects high PL intensity when continuously illuminating the NV center.

If the electron spin state is instead |m_s = ±1>, two possible decay paths are possible. The first one is again spin-preserving and leads to the previously-described PL emission. The second one is nonspin-preserving and the system ends up in the spin state |m_s = 0> of the ground level (see Figure 1a). In this latter case, the system relaxes through the two metastable singlet spin states, namely, $^1A$ and $^1E$. These have a transition wavelength in the infrared range, $\lambda = 1046$ nm. This process is referred as “intersystem-crossing” because it involves a transition between...
Figure 1. a) NV\(^{-}\) state transition that occurs after laser excitation and MW excitation. The coupling of the states \(|m_s = \pm 1\rangle\) with the metastable level generates a statistically lower PL emission than the \(|m_s = 0\rangle\). b) PL collected from the NV\(^{-}\) center as a function of the MW frequency. A dip in correspondence of the zero-field splitting \(D_{gs}\) (resonance frequency of the undisturbed NV\(^{-}\) center, at room temperature) can be observed. c) Sketch of the differential measurement. From the ODMR spectrum (upper part of the figure) the differential spectrum (lower part of the figure) is derived taking, for every value of the microwave (MW) frequency, the difference \(\tilde{F} = F_1 - F_2\) in PL between two points \((F_1, F_2)\) separated by \(2f_{dev}\). \(\tilde{F}\) is zero at the two extremes of the spectrum and at the resonant frequency \(D_{gs}\). Around \(D_{gs}\) there is a region where \(\Delta \tilde{F}\) depends linearly on \(\Delta D_{gs}\) through the differential spectrum slope. d) Example of ODMR and differential ODMR spectra at two different temperatures.

Due to the presence of this nonradiative decay path the spin state \(|m_s = \pm 1\rangle\) is less luminescent with respect to \(|m_s = 0\rangle\) one. Remarkably, the above-explained feature allows the optical read-out of the spin state of the NV center at room temperature. This is achieved by the method of ODMR which consists in applying a microwave field of variable frequency, simultaneously with irradiation of a nonresonant laser able to excite the electronic state. In this case, a 532 nm pump laser is required. When the microwave frequency matches the exact resonant frequency \(D_{gs}\) between the \(|m_s = 0\rangle\) and \(|m_s = \pm 1\rangle\) states, less PL is collected and therefore a fluorescence dip is observed in the ODMR spectrum (see Figure 1b). This allows the optical readout of the spin state.
It is important to underline that when NDs are used as sensors instead of bulk diamonds, the degeneration of the $|m_t = +1\rangle$ and $|m_t = -1\rangle$ levels are removed by mechanical stress. In this case, two dips are observed in the ODMR spectrum. Both of them show a shift in the same direction in the resonance frequency as the temperature varies (Figure 1d).

By applying this protocol, a dip in the collected PL intensity is observed when the microwave frequency matches the exact resonant frequency $D_g$, between the $|m_t = 0\rangle$ and $|m_t = \pm 1\rangle$ electronic spin states of the NV center. Since the value of $D_g$ depends on temperature, this technique allows determining the temperature at the ND location by measuring a frequency shift.

The temperature variation $\Delta T$ measured with the ND sensor is estimated through a direct PL change observation according to the equation

$$\Delta T = \frac{\Delta f}{\text{slope} \cdot \partial D_g / \partial T}$$

(1)

where $\Delta F$ is the actual measured physical signal (difference in the photocounting rate), slope represents the constant of proportionality that connects $\Delta F$ to the resonance frequency shift $\Delta D_g$. Finally, $\partial D_g / \partial T$ represents the coupling constant. Under ambient conditions the coupling constant $\partial D_g / \partial T$ is estimated as $-75$ kHz °C$^{-1}$ for a bulk diamond (see the Experimental Section). The sensitivity of the measurement method is about $3$ °C$^{-1}$, comparable to the state of the art in detection of biological signals with NV sensors in NDs.

2.2. Differential ODMR Spectrum Experimental Measurement Technique

In the differential spectrum measurement technique applied in this experiment, the full ODMR spectrum was acquired only once for 60 s. The data was postprocessed creating the differential spectrum (see Figure 1c). In correspondence to each microwave scanning frequency $f$, a differential photoluminescence (PL) value $\Delta F$ was associated as follows: $\Delta F = F(f) - F(f)_{\text{ref}}$.

Here $F$ represents the PL collected from the ND sensor; $f_{\text{ref}} = f + f_{\text{dev}}$, $f_{\text{dev}} = f - f_{\text{dev}}$, and $f_{\text{dev}}$ is a frequency deviation chosen in order to optimize the differential graph, in our case $f_{\text{dev}} = 2$ MHz.

The method advantage is that in correspondence of temperature variations, for a fixed frequency $f$ the differential signal $\Delta F$ results linearized with respect to frequency shift around the resonant frequency $D_g$.

$$\Delta F = \text{slope} \cdot \Delta D_g$$

(2)

The slope value was then obtained through a linear regression from the differential spectrum.

The differential signal $\Delta F$ was acquired only at one frequency chosen in the linear region of the differential spectrum. The signal $\Delta F$ was acquired in two steps: in the first step the frequency generator was set at $f_1$ for 50 ms; in the second step the frequency was switched to $f_2$ for the same sampling time. Between the two steps a waiting time of 10 ms was set. The cycle was repeated for 500 times (effective repetition rate 8.3 Hz). Any variation $\Delta F$ recorded was attributable to a shift in the resonance frequency $\Delta D_g$ and therefore to a temperature variation $\Delta T$, according to Equation (1).

This differential method is analog to probe the fluorescence at the inflection point of the ODMR spectrum, see Figure 1a. It removes the effects of extrinsic background signal and mitigates the effects of variations in the NV$^-$ fluorescence and in resonance linewidth.\[40\] To further reduce the effect PL stability and the spatial inhomogeneities of the microwave field we use a very stable laser and specially designed planar broadband antenna, which provided a strong, homogeneous, electromagnetic radiation.\[41\]

Compared to four frequency method used,\[71\] our method is of simpler experimental implementation and it is suitable when no signal from magnetic field is expected. The fluctuating magnetic field in our case represents a source of uncertainty.

2.3. Temperature Variation Measurement

The thermometric apparatus used for the experiment was based on a single-photon-sensitive confocal microscope compatible with ODMR measurement. A simplified scheme of the apparatus is shown in Figure 2a. The optical excitation power provided by a CW 532 nm laser (Coherent Promethes 100NE, Relative Intensity Noise $< -90$ dB Hz, 2nd harmonics) was $< 1$ mW and the acquisition time of all PL measurements was 60s. These values were chosen as a trade-off between the fast measurement and the precise temperature estimation. We underline that our choice guarantees cell viability and, in principle, allows for time-resolved measurements on the scale of metabolic processes.

To exploit the ODMR protocol for measuring intracellular temperature variations, hippocampal neurons (10 days in vitro) were incubated for 5 h with 0.6 μg mL$^{-1}$ NDs (average diameter 185 nm, PDI 0.115), which is far below the cytotoxicity threshold ($< 250$ μg mL$^{-1}$), as demonstrated in a previous work of our group,\[26\] and by\[42\] which tested the ND cytotoxicity up to 250 μg mL$^{-1}$. ND internalization (Figure 2c) was assessed by means of standard confocal imaging (see the Experimental Section).

In the following, all measurements were performed after checking that the PL counts were stable and no unwanted movement (drift, displacement, Brownian motion) of the NDs took place.

To detect the temperature variations associated with different patterns of neuronal activity, recordings were performed in three conditions (Figure 2b): i) in the presence of external Tyrode solution (CTRL), since in this condition the hippocampal network was spontaneously firing;\[28\] ii) after addition of 100 μM of picrotoxin, a selective GABA$\alpha_2$ inhibitor, which is known to drastically potentiate the neuronal firing activity;\[31\] iii) after subsequent addition of 0.3 μM tetrodotoxin (TTX) and 500 μM cadmium chloride (Cd), to inhibit the spontaneous firing. Bath temperature was kept at 37 °C and the photocounting rate was performed from internalized NDs. With reference to the above conditions, a more detailed account of the three measurements condition follows hereby.

- **Condition (i):** To exclude that the perfusion system could induce some temperature variations, the photocounting rate $\hat{F}$ from an internalized ND was measured before and after perfusing the cells with Tyrode solution (CTRL, in Figure 2d). From
the photoluminescence difference $\Delta \tilde{F}$, the temperature variation was estimated according to Equation (1). In these conditions we revealed an average temperature variation $\Delta T = (0.05 \pm 0.15) \degree C$ (weighted average evaluated on a sample with a numerosity $N = 11$), proving that no significant temperature variations are associated with thermal exchanges of the perfusion system.

After this preliminary test, the next step was to assess whether changes in neuronal firing could be associated to temperature variations.

- Condition (ii): $\Delta T$ was estimated by comparing $\tilde{F}$ before and after perfusing the cells with a picrotoxin-enriched Tyrode solution. The addition of the GABA$_A$ inhibitor caused a significant temperature increase, $\Delta T = (1.02 \pm 0.24) \degree C$ ($N = 10$), associated with the increased firing rate (Figure 2a). As a more quantitative corroboration of our conclusions, a thorough statistical uncertainty analysis was performed (see the Experimental Section) to confirm that a significant localized temperature variation occurs in a neuron deriving from sustained firing activity ($1.12 \times 10^{-8}$ significance in a Welch $t$-student test).
Figure 3. Laser exposure does not affect the hippocampal neurons. a) modulation of the firing activity by PICRO (MEA recordings): representative traces from the same electrode under control condition, + PICRO. Insets: higher magnification of single spikes and bursts. b) Histogram of mean frequency in the different experimental conditions. The statistical difference of PICRO respect to other conditions is indicated ($p < 0.05$, *). c) Left: Representative traces in control condition and after laser irradiation and (right) corresponding raster plot. In the raster plot, 7 representative channels (ch_1÷ ch_7) are shown.

- Condition (iii): network activity was silenced by means of the Tyrode solution enriched with TTX+Cd. In this case, $\Delta T$ was measured by comparing $\bar{F}$ before and after perfusing the cells with picrotoxin + TTX+Cd solution. Under this last condition, a significant temperature decrease $\Delta T = (−0.50 \pm 0.17) ^\circ C$ ($N = 8$) was revealed (Figure 2d) with a $2.21 \times 10^{-6}$ significance in a Welch t-student test (see the Experimental Section). These data demonstrate that both potentiation and silencing of the neuronal network activity could be assessed by temperature variations and that the hippocampal network, exhibiting basal spontaneous activity, displays higher temperatures than the completely silenced network.

Another set of experiments was carried out for measuring the average temperature variation associated to noninternalized ND before and after perfusing the cells with picrotoxin solution: in this case no significant difference with respect to the control conditions was found: $\Delta T = (−0.04 \pm 0.23) ^\circ C$ ($N = 7$). This crossover trial excluded that temperature variation can be due to other external factors than the drug-potentiated neuronal activity.

Furthermore, it has been tested whether the laser exposure, microwave radiation and incubation with NDs could affect the spontaneous firing activity of hippocampal neurons during ODMR measurement protocol (as detailed in the Experimental Section). As shown in Figure 3, both the spontaneous firing rate and the action potential waveform were not significantly altered by laser and ND exposure, confirming that cell excitability and ion channel functioning are preserved, in good agreement with our previous findings.[26]

Finally, in parallel with the thermometry experiment, the effect of GABA$_A$ receptor blockade on burst-firing of hippocampal neurons was assessed by comparing the spontaneous burst ratio of cells responses before and after application of picrotoxin. These recordings have been performed by means of conventional MEAs (MCS, multichannel system). An expected significant increase of the spontaneous activity,[44] (from $2.1 \pm 0.3$ to $3.4 \pm 0.3$ Hz), which is associated to disinhibition of GABA$_A$ receptors (Figure 3a), and a complete silencing of the network in the presence of the Na$^+$ and Ca$^{2+}$ voltage-dependent channel blockers TTX and Cd was observed. The data presented allowed to conclude that the
increase in temperature recorded in the experiment is directly related to the altered firing activity of hippocampal neurons.

3. Conclusion

Intracellular localized temperature gradients are associated with metabolic activity and with a variety of reactions and intracellular processes.

Real-time mapping of intracellular subtle temperature gradients and event-driven increased temperature represents a tool of utmost interest for monitoring the functional activity of the cells, for identifying localized signaling under physiological conditions and for exploitation as a diagnostic tool.

Here we applied the ODMR technique to cultured hippocampal neurons for monitoring temperature variations in different conditions of network excitability. We provided evidence that a significant temperature increase can be correlated with the altered firing activity of cultured hippocampal neurons, in our case induced by picrotoxin, and that this phenomenon can be successfully observed at the single-cell level exploiting nanosensor based on NDs.

It is worth noticing that picrotoxin-induced disinhibition primarily alters cell excitability, but the increased action potential firing activity may in turn alter intracellular pathways: since we cannot predict whether NDs are localized at the plasma membrane or at the intracellular level, the detected temperature increase during perfusion with picrotoxin is ascribed to sustained firing activity and/or potentiated cell metabolism.

The measured temperature increases can be used to reveal the onset of different intracellular processes other than membrane-delimited pathways, likely involving an altered cell metabolism as here was demonstrated with picrotoxin. Thus, our future goals will be devoted to functionalizing NDs to subcellular compartments, to detect thermogenesis at specific subcellular sites.

These results, backed by a thorough analysis of other possible alternative causes of temperature variation, pave the way to a systematic study of cell activity with impacts ranging from a better insight on current unknowns associated with cells functioning, such as, e.g., the aforementioned discrepancy between experimental data and thermal diffusion models, to the research on specific pathologies.

Our findings also prompt further applications where ongoing advances in microelectrode array technology (MEA), combined with quantum sensing paves the way for experiments that take advantage from the synergy of the two techniques, such as synchronized measurements of cellular activity with metabolism processes or propagation of electromagnetic signals.

4. Experimental Section

Experimental Setup: The thermometric apparatus used for this experiment (Figure 4) was based on an Olympus IX73 inverted microscope, in which optical elements to implement single-photon confocal imaging and microwave control for ODMR measurements were integrated. A CW 532 nm laser (Coherent Prometheus 100NE, noise reduced regime, 2nd harmonics), attenuated down to 1 mW, was used to excite the NV centers in the NDs. Then an acousto-optic modulator (AOM), controlled by a pulse generator (Pulsar Generator, Pulse Generator, Swabian Instruments), was applied to the laser emission to have the cell irradiated only during the measurement. Finally, a 60x air microscope objective (Olympus UPLANFL, NA = 0.67) was used for both excitation and photoluminescence (PL) collection. The spot size of the focused laser beam was the 1.2 × 1.3 μm². The NV PL was filtered by a 567 nm dichroic mirror, a 650 nm long-pass filter and a Notch filter centered at 532 nm to remove the residual green laser scattering, and then collected by a single photon avalanche diode (SPAD, SPCM-AQR 15, Perkin Elmer). The emission rate from a single ND was around 300 counts s⁻¹.

The Petri dish, containing the cell culture with NDs, was placed in a closed illumination chamber with a temperature control (Okolab Temperature Controller, temperature stability 0.1 °C). The sample could be moved via a manually or software-controlled piezoelectric XYZ scanning stage.
Allowing selecting the area of interest. The output signal from the SPAD was fed to a data acquisition board system (National Instruments, USB-6343 BNC).

To implement the ODMR measurements, a microwave source (Keysight N5172B) was used. The output signal was then amplified (Mini-Circuits, ZHL 16W 43+) to a power of 20 dBm and fed to a homemade planar broadband antenna, which provided a strong, homogeneous, electromagnetic radiation. The Petri dish was placed on the top of the antenna.

Experimental Procedure: The experimental procedure for detecting temperature variation using NV centers in NDs was the following.

A Petri dish containing cultured hippocampal neurons (10 DIV), previously exposed to NDs, was positioned inside the single-photon confocal microscope incubation chamber. Using the motorized stage, the sample was scanned by incident laser (P < 1 mW, RIN < −90 dB Hz⁻¹) in order to select a single ND inside a cell. The ODMR spectrum was acquired for 60 s and the data were postprocessed to create the differential. Then, the microwave frequency was set in accordance with the minimum of the ODMR spectrum (zero of the differential spectrum). The differential signal of luminescence at this frequency was acquired for 60 s. The microwave field before and after data acquisition was always turned on to avoid possible heating transient processes.

In a successive step, solutions containing 100 μm picotoxicin and then 300 nM TTX + 200 μm CdCl₂ were added to the medium. The measurement of the differential signal was repeated at the same resonant frequency as in the previous step. Finally, the ODMR spectrum was reacquired for 400 s, recreating a new differential spectrum in order to improve the statistics and to evaluate the linear region slope with lower uncertainty. The temperature variation ΔT was evaluated through a direct PL change measurement according to equation

\[ \Delta \tilde{F} = \text{slope} \cdot \partial D_{gs} / \partial T \cdot \Delta T \]  

Here \( \Delta F \) represents the differential signal variation (difference in the photon counting rates), evaluated at the linear region of the differential spectrum and associated with a temperature variation \( \Delta T \). The slope represents the constant of proportionality that connects \( \Delta \tilde{F} \) to the resonance frequency shift \( \Delta D_{gs} \). Finally, \( \partial D_{gs} / \partial T \) represents the coupling constant.

**Experimental Procedure—Estimation of \( \partial D_{gs} / \partial T \):** Under ambient conditions the coupling constant \( \partial D_{gs} / \partial T \) is estimated as 75 kHz °C⁻¹ for a bulk diamond [53,54]. For NDs this constant varies due to the different strain splitting and, in principle, should be calibrated.

The method to obtain \( \partial D_{gs} / \partial T \) was the same adopted for the cell temperature measurement, with the difference that in this case \( \Delta T \) was not induced by a substance perfusion but forced by the incubator chamber’s heater, and that \( \Delta T \) is measured by a thermocouple. The procedure was the following. The NDs were placed on the Petri dish with distilled water. The ODMR spectrum was acquired and so was the differential spectrum by postprocessing the data, see above. The linear region slope of the differential spectrum was obtained by performing a linear regression. The microwave generator was set at the resonant frequency \( D_{gs} \) that occurred in the initial temperature condition reached in the incubator chamber, at which the differential spectrum signal \( \tilde{F} \) was zero. The incubator temperature was varied, recording it with the thermocouple. The new value of the differential signal was read. Finally, \( \Delta F \) was plotted as a function of the temperature increase \( \Delta T \) recorded by the thermocouple. Through a linear regression the constant of proportionality \( b \) was evaluated, i.e., the quantities mentioned obey the following physical relationship

\[ \Delta \tilde{F} = b \cdot \Delta T \]  

where

\[ b = \text{slope} \cdot \partial D_{gs} / \partial T \]  

By inverting the last equation, the quantity \( \partial D_{gs} / \partial T \) was estimated. The results are shown in Table 1. Both the arithmetic mean and the weighted mean coincide: \( \partial D_{gs} / \partial T = (-76 ± 4) \) kHz °C⁻¹. We used this estimated value for \( \partial D_{gs} / \partial T \) and the associated uncertainty in the data analysis.

**Cell Preparation and Nanodiamonds (ND) Labeling:** Hippocampal neurons were obtained from C57BL/6 16-day embryos. All animals were housed under a 12 h light/dark cycle in an environmentally controlled room with food and water ad libitum. All experiments were conducted in accordance with the European Community’s Council Directive 2010/63/UE and approved by the Italian Ministry of Health and the Local Organism responsible for animal welfare at the University of Turin (Authorization 695/2020-PR). Hippocampus was rapidly dissected under sterile conditions, kept in cold HBSS (4 °C) with high glucose, and then digested with papain (0.5 mg mL⁻¹) dissolved in HBSS plus DNase (0.1 mg mL⁻¹). Isolated cells were plated at density of 1200 cells mm⁻² onto the MEA and 1000 cells mm⁻² onto the glass Petri dishes. Both the MEAs and the dishes were previously coated with poly-DL-lysine and laminin, this allowed the neurons to adhere to the center of the device by using a ring made of Sylgard 184 (Dow Corning), which was removed after 4 h. The cell medium is composed of 1% penicillin/streptomycin, 1% glutamax, 2.5% fetal bovine serum, 2% B-27 supplemented neurobasal medium. The neurons were

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<th>$\frac{\partial D_{gs}}{\partial T}$</th>
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incubated in a humidified 5% CO₂ atmosphere at 37 °C. Recordings were carried out at 10–12 days in Vitro (DIV).

The initial ND solution (1 mg mL⁻¹, MilliQ water) was sonicated for 5–10 min (100 W power, 80 kHz frequency), in order to separate diamond particle clusters into single NDs. After sonication, the 60 μL of solution was diluted with 1 mL of Tyrode to obtain a final ND concentration of 0.6 μg mL⁻¹ which was poured into the Petri dish. The dish was exposed to NDs for 5 h in order to allow NDs internalization. Once the dish was extracted from the incubator, the cell medium was removed and replaced with 2 mL of Tyrode solution.

Micro Electrode Array Recordings: Multisite extracellular recordings were performed using the MEA-system, purchased from Multi-Channel Systems (Reutlingen, Germany). The 60 electrodes array (TiN) was composed of an 8 x 8 square grid with 200 μm inter-electrode spacing and 30 μm electrode diameter. Data acquisition was controlled through MC_Rack software (Multi-Channel Systems Reutlingen, Germany), sampling at 10 kHz. Experiments were performed in a non-humidified incubator at 37 °C and with 5% CO₂, maintaining the culture medium. Before starting the experiments, cells were allowed to stabilize in the non-humidified incubator for 5 min; recording of the spontaneous activity was then carried out for 120 s. Mean frequency has been evaluated over 120 s recording. The data are analyzed using Neuroexplorer software (Nex Technologies, Littleton, MA) after spike sorting operations.

Confocal Image Acquisition: After 5 h exposure to ND (0.6 μg mL⁻¹), the cytoplasmic labelling dye (CellTracker Green CMFDA, Thermafischer) was added to the medium. This allowed to identify the cell boundaries of cultured hippocampal neurons (10 DIV) together with the internalized NDs, characterized by a red fluorescent emission. For standard confocal imaging, hippocampal neurons were plated on 35 mm dishes (ibidi GmbH, Planegg/Martinsried, Germany). Cells were analyzed using a confocal laser scanning microscope Leica TCS SPS, equipped with an argon ion and a 561 nm DPSS laser. Cells were imaged using a HCX PL APO 63x/1.4 NA oil immersion objective at a pixel resolution of 0.08 × 0.08 × 0.3 μm. The luminescent emission from the fluorescent NDs was excited by 561 nm laser, while the emission was collected in the 650–750 nm spectral range. The same excitation wavelength was ineffective in untreated neurons. Green fluorescence for intracellular staining was obtained by 488 nm wavelength. Image analysis was performed using ImageJ software.

Monodisperse ND Preparation: NDs were supplied by Microdiamant Switzerland (MSY 0-0.25, containing ≈100–200 ppm of natural nitrogen impurities) and treated as described previously. Briefly, the NDs were oxidized by air oxidation at 510 °C for 5 h and then wet oxidized in a HF:HNO₃ 2:1 v/v stirred mixture at 160 °C for 2 days in a PTFE container. The acids were removed using consecutive centrifugation/washing and the resulting pure oxidized NDs were lyophilized. The monodisperse NDs (hydrodynamic diameter 205 nm) were isolated using differential centrifugation in water based on gradual centrifugation/dilution steps, combination of the separated pellets, followed by their lyophilization.

Fluorescent NDs Irradiation and Oxidation: A total of 330 mg of the monodisperse ND was irradiated at 870 °C in an external target for 80 h with a 15.7 MeV electron beam (2.5 × 10¹⁰ particles cm⁻²) extracted from the MT-25 microtron. After irradiation, the NDs (315 mg) were annealed at 900 °C for 1 h under argon atmosphere and subsequently oxidized in air for 5 h at 510 °C. The resulting fluorescent NDs were wet oxidized using a 2:1 v/v mixture of HF:HNO₃, washed and lyophilized, providing 212 mg of NDs with NV centers (yield 64%). Before use, the lyophilized NDs were redispersed in MilliQ water using cup horn sonication to concentration 1.0 mg mL⁻¹.

Dynamic Light Scattering: The hydrodynamic particle diameter was measured using a Zetasizer Nano ZS (Malvern Instruments) in disposable transparent cuvettes. The samples were measured at Room Temperature with a concentration of 0.05 mg mL⁻¹ and a total volume of 1 mL.

Figure 5. Validation of temperature detection by ND sensor. The legend box shows the temperature values recorded by the thermocouple. The box plot shows the temperature values recorded by the ND sensor. The mean and its uncertainty (standard deviation divided by the square root of repeated measurements) are reported as horizontal lines in each data set (N = 10, 60 s acquisition). The starting point is highlighted as an asterisk. The incubator temperature at the end of the cycle is consistent with the initial temperature within the incubator stability (0.1 °C).
Table 2. Summary of temperature variation $\Delta T \pm \sigma$ ($^\circ C$) for the three (sample size, respectively, $N = 11, N = 10, N = 8$) independent groups.

<table>
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<th>CTRL</th>
<th>+PICRO</th>
<th>+PICRO+TTX+Cd</th>
</tr>
</thead>
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<tr>
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<td>$0.14 \pm 0.48$</td>
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<td>$-0.17 \pm 0.47$</td>
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<tr>
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<tr>
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All the statistical analysis was performed by using Wolfram Mathematica software.

Acknowledgements

The authors acknowledge Prof. E. Carbone for useful discussion and C. Franchino for cell preparation. This project (No. 20IN0605) QADeT leading to this publication has received funding from the EMPIR program co-financed by the Participating States and from the European Union’s Horizon 2020 research and innovation program. The project was further supported by PATHOS EU H2020 FET-OPEN Grant No. 823946 Horizon 2020, and European Regional Development Fund; OP RDE; Project CARAT (No. CZ.02.1.01/0.0/0.0/16_026/000382), Compagnia di San Paolo, PROGETTO TRAPEZIO.

Conflict of Interest

The authors declare no conflict of interest.

Author Contributions

G.P. and G.T. contributed equally to this work. M.G., V.C., E.B., E.M., I.P.D., P.C., and P.T. planned the experiment. The main experimental work was equally realized by G.T. (for the biological part) and G.P. (optical one) with the direct help and supervision of E.B. and E.M. (for the optical one) in the laboratories directed by V.C. (biological one) and M.G. (optical one). The quantum sensing facility at INRIM was developed by E.M., E.B., and P.T. All the authors equally contributed to the discussion on experimental settings and on data analysis, as well as in writing the paper. P.C. and K.K. provided the nanodiamond samples.

Data Availability Statement

The data that support the findings of this study are available from the corresponding author upon reasonable request.

Keywords

intracellular nanoscale sensing, nanodiamonds, nitrogen-vacancy (NV) centers, ODMR

Received: May 30, 2022
Revised: June 28, 2022
Published online: July 25, 2022
