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Linear response of one-dimensional liquid ^4He to external perturbations

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Abstract We study the response of one-dimensional liquid ^4He to weak perturbations relying on the dynamical structure factor, $S(q, \omega)$, recently obtained via *ab-initio* techniques [*Phys. Rev. Lett.* **116**, 135302 (2016)]. We evaluate the drag force, F_v , experienced by an impurity moving along the system with velocity v and the static response function, $\chi(q)$, describing the density modulations induced by a periodic perturbation with wave vector q .

Keywords Luttinger liquid · one-dimensional superfluidity · helium · drag force · static density response function

1 Introduction

One-dimensional (1D) systems occupy a unique place in the realm of many-body quantum mechanics, due in particular to the loosened distinction between the behavior of fermions and bosons with hard-core repulsive interaction. Electronic transport properties in quasi-1D have been extensively investigated in various systems such as fractional quantum Hall edge states [1], carbon nanotubes [2,3], conducting conjugated polymer nanowires [4], and semiconductors nanowires [5]. Also quasi-1D bosonic systems realized with ultracold atoms confined in cigar-shaped traps have attracted considerable attention [6]. 1D behavior has also been theoretically predicted and experimentally observed to play a crucial role for the characterization of ^4He atoms moving inside dislocation lines in crystalline helium samples [7,8,9], or confined inside nanopores [10,11,12,13,14,15].

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The low-energy properties of a vast class of 1D systems, including liquid ^4He , are well captured by the Tomonaga-Luttinger liquid theory (TLL) [16, 17, 18]. The TLL theory is an effective low-energy field theory governed by an exactly solvable quadratic Hamiltonian that depends only on two parameters, namely the sound velocity c and Luttinger parameter K_L , which characterize the large-distance behavior of the correlation functions and the thermodynamic properties of the system. For spinless Galilean-invariant systems, as the ones we are going to consider here, these two parameters are related via $c = v_F/K_L$, where $v_F = \hbar\pi\rho/m$ is the Fermi velocity of a 1D ideal Fermi gas with the same density $\rho = L/N$ of the system. Thus the properties of the system depends only on a single parameter, K_L , that turns out to be proportional to the square root of the compressibility κ_S : $K_L = \sqrt{\hbar^2\pi^2\rho^3\kappa_S/m}$.

Since the Mermin-Wagner theorem [19] prevents the existence of Bose-Einstein condensation in 1D systems with short-range velocity-independent interaction, the standard picture of superfluidity relying on the order parameter provided by the condensate wavefunction [20] has to be generalized [21]. However, in the TLL phase, correlation functions feature a power-law decay [18] and the system is superfluid in the sense that it displays a quasi-off-diagonal long-range order [6, 21].

In this paper we characterize the dynamical properties of 1D liquid ^4He by computing the friction or drag force, F_v , exerted on a particle of mass m moving with velocity v inside the system, which provides a generalization of Landau's celebrated criterion of superfluidity [22, 23]. We determine also the static density response function, $\chi(q)$, which is of paramount importance for the response of the system to a periodic potential, relevant for the interpretation of experiments [7, 24, 25] and for density functionals theories [26, 27, 28, 29].

Both F_v and $\chi(q)$ are obtainable from the knowledge of the dynamical structure factor, $S(q, \omega)$, that we have recently achieved via an *ab-initio* full microscopic approach [30], by combining the Genetic Inversion via Falsification of Theories (GIFT) algorithm [31, 32, 33, 34, 35, 36, 37, 38] to perform the analytic continuation of imaginary-time correlation functions computed via the exact zero-temperature Path Integral Ground State (PIGS) method [39, 40, 41].

2 Drag Force

According to Landau [22, 23], an obstacle in a fluid, moving with velocity v , may cause transitions from the ground state to excited states lying on the line $\epsilon(k) = \hbar v k$ in the energy-momentum space. If all the spectrum is above this line, the motion cannot excite the system, and the flow of the impurity is frictionless. However, even when the line intersects the spectrum, the transition probabilities to these states can be strongly suppressed due to the interaction or to the external perturbing potential. In this case, the drag force provides a quantitative measure of superfluidity.

The drag force, F_v , exerted on an impurity moving with velocity v in a 1D medium, can be computed relying on Fermi's golden rule, by expressing the energy loss per unit time due to the impurity, $\frac{d\mathcal{E}}{dt} = -F_v v$, as [42]:

$$\frac{d\mathcal{E}}{dt} = - \sum_{q(=k_f-k_i)} p(k_i \rightarrow k_f) \hbar\omega_q = - \int dq |V_q|^2 \rho \hbar\omega_q S(q, \omega_q) \quad (1)$$

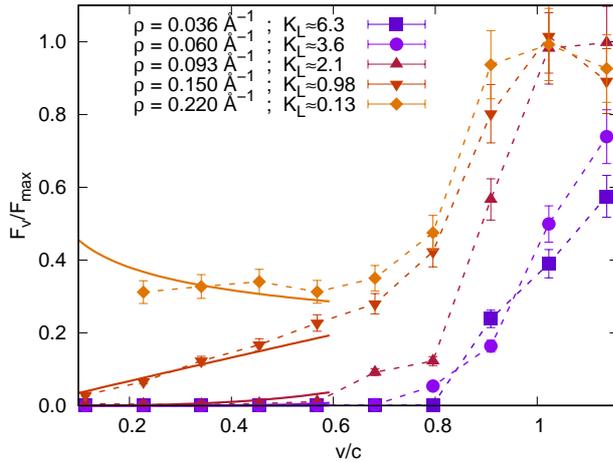


Fig. 1 Drag force F_v as a function of the ratio v/c , compared to the Luttinger liquid theory prediction for small v (solid lines) and normalized to the maximum value F_{max} . Corresponding values of the Luttinger parameter are also indicated. Errorbars are estimated from the typical variance of the reconstructed GIFT spectra. (Color figure online).

where $p(k_i \rightarrow k_f)$ is the rate for the scattering process, $\omega_q = \hbar(k_f^2 - k_i^2)/2m = qv - \hbar q^2/2m$ and V_q the Fourier transform of the impurity-medium interaction V . For a heavy impurity, $v \gg \hbar q/m$, (1) simplifies to:

$$\frac{d\mathcal{E}}{dt} = - \left(\int_0^\infty dq |V_q|^2 \rho \hbar q S(q, qv) \right) v = -F_v v \quad (2)$$

Equation (2) shows that knowledge of $S(q, \omega)$ and of V_q grants the ability of computing the drag force. The relationship between the concept of drag force and Landau's criterion for superfluidity is readily understood: if $S(q, \omega)$ is concentrated above the straight line $\hbar v_0 q$ in the momentum-energy plane, then $F_v \equiv 0$ for all velocities $v < v_0$. On the other hand, the integral can be very small or vanish even if the spectrum lies below the line $\hbar v_0 q$, but the excitation probabilities are suppressed [42], for instance because V_q takes non-zero values only in a finite region of momentum space. In a broad class of 1D systems, like the ideal Fermi gas, ^4He at high density [30] and hard rods [43], $S(q, \omega)$ touches the $\omega = 0$ line with finite weight at $2k_F$ ($k_F = \pi\rho$ being the Fermi wave-vector), whence the possibility of a dissipationless flow is by no means obvious.

To compute the drag force (2), we have approximated V_q by using a contact potential with the same scattering length of the repulsive part of the ^4He interaction (as in Ref. [44]), and the $S(q, \omega)$ estimated in Ref. [30]. Our results are shown in Fig. 1. We observe an increase of F_v with the velocity, and a saturation to a maximum value F_{max} at $v \simeq c$. The increase is consistent with Luttinger liquid theory [42, 45]. In fact, for a slow impurity $v \ll c$, the most important contribution to (2) comes from the region $q \simeq 2k_F$. As shown in Ref. [45], for $q \rightarrow 2k_F$ and $\omega \rightarrow 0$, $S(q, \omega)$ has the power-law behavior:

$$S(q, \omega) \propto \omega^{2(K_L-1)} (1-x^2)^{K_L-1} \quad (3)$$

being $x = c(q - 2k_F)/\omega$, with the constraint $|x| \leq 1$. Inserting (3) into (2) yields the following expression for the drag force:

$$F_v \propto \int_{-1}^{\min(\frac{c}{v}, 1)} \frac{dx}{(1 - \frac{v}{c}x)^3} \left(\frac{\frac{v}{c}}{1 - \frac{v}{c}x} \right)^{2(K_L-1)} (1 - x^2)^{K_L-1}. \quad (4)$$

For $v \ll c$, the approximation $1 - \frac{v}{c}x \simeq 1$ is accurate and one finds $F_v = F_0(v/c)^{2K_L-1}$ [42, 45], that compares satisfactory with our estimates of F_v .

We remark that the drag force is non-vanishing at any v , but for $K_L > 1/2$ its power-law behavior determines a superfluid response of the system, allowing for impurities to flow with small dissipation [46]. Notice also that, for $K_L < 1/2$ (e.g. at the density $\rho = 0.220 \text{ \AA}^{-1}$), the previous analytical expression implies a divergence of F_v for small velocities; nonetheless, the physical dissipated power in Eq.(2) is always vanishing at small velocity, since it behaves as $\dot{E} \propto v^{2K_L}$. Our estimate of F_v is perturbative, therefore it is relevant for soft impurities such as small geometry deformations in quasi-1D systems. It is clear that for hard-core impurities (like ^3He pinned to ^4He dislocations [47, 48]) the superfluid response is completely suppressed.

3 Static density response function

The static density response function characterizes the effect of a static periodic perturbation on a homogeneous system [49, 50, 51, 52], and can be computed relying on the Hellman-Feynman theorem [51] or from the first negative moment, $m_{-1}(q)$, of $S(q, \omega)$:

$$\chi(q) = -\frac{2\rho}{\hbar} \int_0^\infty d\omega \frac{S(q, \omega)}{\omega} = -\frac{2\rho}{\hbar} m_{-1}(q). \quad (5)$$

Our results are shown in Fig. 2.

In the low-momentum regime, according to the TLL theory the dynamical structure factor is concentrated around the dispersion relation $\omega_{LL}(q) = c|q|$, and the Feynman approximation $S(q, \omega) = S(q) \delta(\omega - \omega_{LL}(q))$ can be safely assumed, with $S(q) \simeq K_L \frac{q}{2k_F}$. Therefore:

$$m_{-1}(q) = \frac{1}{E_F} \frac{S(q)^2}{\left(\frac{q}{k_F}\right)^2} \simeq \frac{1}{E_F} \frac{K_L^2}{4} \quad (6)$$

being $E_F = \hbar^2 k_F^2 / 2m$ the Fermi energy. This relation is equivalent to the well known compressibility sum rule obeyed by $\chi(q)$ in higher dimensions [53]. The resulting values for the different considered densities are reported as solid lines in Fig. 2, instead of the obtained $\chi(q)$, that at low momenta is plagued by frequency-discretization effects unavoidable with the actual implementation of GIFT.

As clearly visible in Fig. 2, $\chi(q)$ features peaks at momenta which are integer multiples of $2k_F$, at the densities $\rho = 0.22 \text{ \AA}^{-1}$ and $\rho = 0.30 \text{ \AA}^{-1}$, where K_L is found to be $K_L = 0.389(2)$ and $K_L = 0.1255(5)$ respectively [30]. The presence of such peaks can be easily justified in the light of the TLL theory, with the following heuristic argument. Assuming that $S(q, \omega)$ has the low-frequency power-law divergence $S(q = 2jk_F, \omega) \propto \omega^{2(j^2 k_L - 1)}$ at $q = 2jk_F$ [30, 45, 54, 55], and has

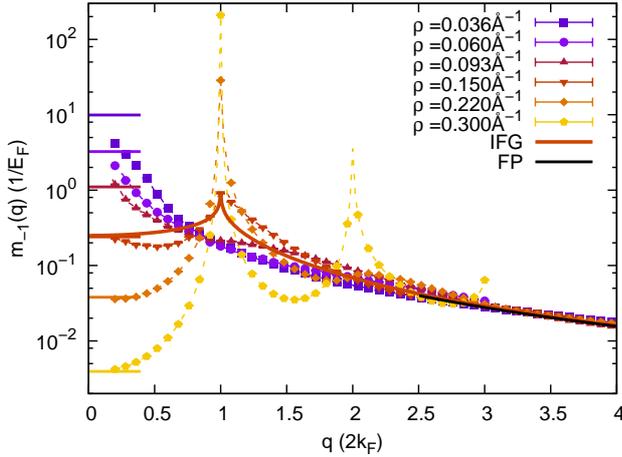


Fig. 2 First negative momentum $m_{-1}(q)$ in units of the inverse Fermi energy $1/E_F$, as function of q in units of $2k_F$. Small horizontal lines at $q \rightarrow 0$ correspond to the asymptotic estimate based on the Feynman approximation. The static response function for the ideal Fermi gas (IFG) and the free particle limit (FP) are also reported as solid lines. (Color figure online).

support above the energy $\hbar\omega_{sc,j} = 4E_F j^2/N$ of the supercurrent state obtained applying a boost of momentum $2jk_F$ to the ground state [42], one can easily prove [56] the presence of peaks in $S(q = 2jk_F)$:

$$S(q = 2jk_F) = \int_{\omega_{sc,j}}^{\infty} d\omega S(q, \omega) \propto \omega_{sc,j}^{2(j^2 k_L - 1) + 1} \propto N^{1 - 2j^2 k_L} \quad (7)$$

because $N^{1 - 2j^2 k_L}$ diverges for $K_L < \frac{1}{2j^2}$ in the thermodynamic limit. Applying the same argument to $m_{-1}(q)$ shows that:

$$m_{-1}(q = 2jk_F) = \int_{\omega_{sc,j}}^{\infty} d\omega \frac{S(q, \omega)}{\omega} \propto \omega_{sc,j}^{2(j^2 k_L - 1)} \propto N^{2 - 2j^2 k_L} \quad (8)$$

diverging in the thermodynamic limit provided that $K_L < \frac{1}{j^2}$. Equation (8) explains the presence of peaks in $\chi(2k_F)$ at the density $\rho = 0.220 \text{ \AA}^{-1}$, and in $\chi(2k_F), \chi(4k_F)$ at $\rho = 0.300 \text{ \AA}^{-1}$.

The $S(q, \omega)$ of 1D Bose systems with repulsive interaction manifests a pseudo particle-hole, typical of Fermi systems, due to a peculiar interplay between dimensionality and repulsion [30]. This suggests that also $\chi(q)$ should display a behavior similar to the static density response function of the ideal Fermi gas, known also as static Lindhard function [57], that we report in Fig. 2 as a solid line. In the large q limit, our results collapse on the asymptotic behavior $m_{-1}(q) \simeq \frac{k_F^2}{q^2}$ as expected also in higher dimensions (free particle limit) [53].

4 Conclusions

We addressed the calculation of linear response functions of one-dimensional ^4He , relying on state-of-art Quantum Monte Carlo calculations of the dynamical struc-

ture factor of the system. In particular, we estimated the force experienced by an impurity moving through the sample and weakly interacting with the helium atoms. This calculation is very interesting since it addresses the fate of the celebrated Landau criterion for superfluidity in a scenario in which the Mermin-Wagner theorem rules out the possibility of Bose Einstein condensation. Our results display a power-law behavior of the drag force as a function of the velocity of the impurity, showing thus a quasi-superfluid response of the system, consistently with the Luttinger liquid paradigm.

We also computed the static density response function $\chi(q)$ of the system, providing the linear response of the helium atoms to an external periodic potential. For interacting quantum systems, this property is notoriously hard to access, requiring either the estimation of dynamical properties or the introduction of an external potential [58]. We suggest interpretation of our data relying on known properties of the dynamical structure factor $S(q, \omega)$, such as the compressibility sum rule in the $q \rightarrow 0$ limit and the free-particle $q \rightarrow \infty$ limit. Varying the density, we observe a crossover from a dilute regime with a smooth static density response function to a quasi-solid high-density regime, where $\chi(q)$ displays peaks whose height increases with the size of the system.

Although our results are valid for a strictly 1D configuration, while experiments with ^4He atoms deal with quasi-1D geometries, such as nanopores or dislocation lines, we hope our findings can provide useful insight and a firm limiting case for realistic situations. In particular, helium in nanopores has been observed to manifest anomalous heat capacity [10] and superfluid response in torsional oscillator experiments [59], with respect to the 2D behavior which is expected from the adsorbed films on the nanopore walls. Extinction of superfluid effects in the limiting 1D case (very small nanopore radius) is elucidated in the present work, depending on the linear density.

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